PARTICLE MOTION THROUGH A DYNAMICALLY DISORDERED MEDIUM: THE EFFECTS OF BOND CORRELATION AND APPLICATION TO POLYMER SOLID ELECTROLYTES

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To study the diffusion of small particles through a dynamically disordered medium, a dynamic bond percolation model has been developed. This differs from the standard percolation theory, in that the lattice is no longer static but undergoes rearrangements which reassign the open and closed bonds. Physically, these rearrangements correspond to orientational motions of the (polymer) host lattice.

An interesting feature of this model is that, even below the percolation threshold, diffusive behavior can occur, as long as the renewal time τ_{ren} , the time characteristic of the rearrangement of bonds, is short compared to the observation time.

Ionic conductivity in polymeric electrolytes is one of the systems for which this model is useful. In these materials, carrier ions diffuse through a medium (the polymer) which is undergoing dynamic motion caused by configurational motions of the polymer. Since polymer chain motions will affect several ion binding sites simultaneously or serially, the model is further elaborated to account for correlations in the segmental motions of the polymer host. The renewal of the bonds is changed from occurring randomly to include simple correlation effects.

Of principal concern in this study is the effect that correlated renewals have on the transport behavior of the model. Simulations were done on a 1-D lattice and a diffusion coefficient calculated. The values of the diffusion coefficients from the two systems (with and without correlated renewal) are studied and their behavior as a function of the fraction of available bonds, f, and the renewal time is compared.

For both correlated and uncorrelated renewals, the systems were diffusive. The diffusion coefficients, in both cases, increased with increasing f and decreasing $\tau_{\rm ren}$, corresponding to an increase in the free volume, the configurational entropy, and the temperature of the polymer systems. The diffusion coefficient from the correlated systems were always smaller than those from the uncorrelated systems, except for the limit f = 100% and $\tau_{\rm ren} << \tau_{\rm hop}$. The ratio of the diffusion coefficients for the correlated and uncorrelated systems was studied as a function of $\tau_{\rm ren}$ and f. This ratio falls off to a constant value as $\tau_{\rm ren}$ is increased and reaches minimum value at f = 50%. This behavior of the ratio as a function of f can be explained by considering the diffusion of the bonds in the lattice for the correlated case.

1. INTRODUCTION

Percolation models are often used to study transport properties in rigid media. In bond percolation schemes, a particle is allowed to hop from site to site on a lattice in which a given fraction of intersite bonds is randomly blocked. This model has been applied to the study of transport in framework solid electrolytes. The development of solvent free polymer-salt complexes as solid electrolytes 3,4 has stimulated interest in the

transport properties of these non-framework ionic conductors. The importance of polymer backbone fluctuations to the transport of the ions has been demonstrated in a variety of systems. 5,6,7 Theories based on free volume and configurational entropy have been used to describe the temperature dependent conductivity of polymeric solid electrolytes. 4,8 Amorphous regions of the polymer host, above its glass transition temperature, have a substantial amount of free volume or configurational

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entropy which increases with temperature and facilitates ion transport.

The application of the bond percolation model to polymer solid electrolytes is possible if the model is modified to account for the segmental motions of the polymer host. In the Dynamic Bond Percolation (DBP) model, the dynamics of the polymer host are taken into account by considering lattice rearrangements which reassign the blocked and unblocked bonds. The application of DBP to polymer electrolytes, including both its relationship to the free volume model, and formal properties of the model itself 10,11,12 have been studied previously. A slightly different model has been recently investigated by Harrison and Zwanzig. 12

The objective of the present study is to alter the DBP model to include correlations among the lattice bonds undergoing rearrangement, and to make a qualitative investigation of how correlation affects the diffusion coefficients obtained from numerical simulations. Work by Kutner and Kehr¹³ has considered a model which formally resembles DBP with correlated renewal. This work looks at diffusion with two kinds of particles where one kind may, in our case, take the part of the migrating bonds. Here we consider a model in which the lattice motion is far more strongly correlated than in refs. 9-11 or 12. It involves permitting open, or available, bonds to move only between one bond site to the next bond site, and may correspond to a polymer electrolyte in which backbone motions indeed move the available ion motion pathways down the lattice. It also characterizes well the one-dimensional case of light particles (the ions) moving in concert with heavy particles (the lattice). and should be useful in describing the coupled rotation and hopping processes which occur in protonic electrolytes.14

2. MODEL AND DISCUSSION

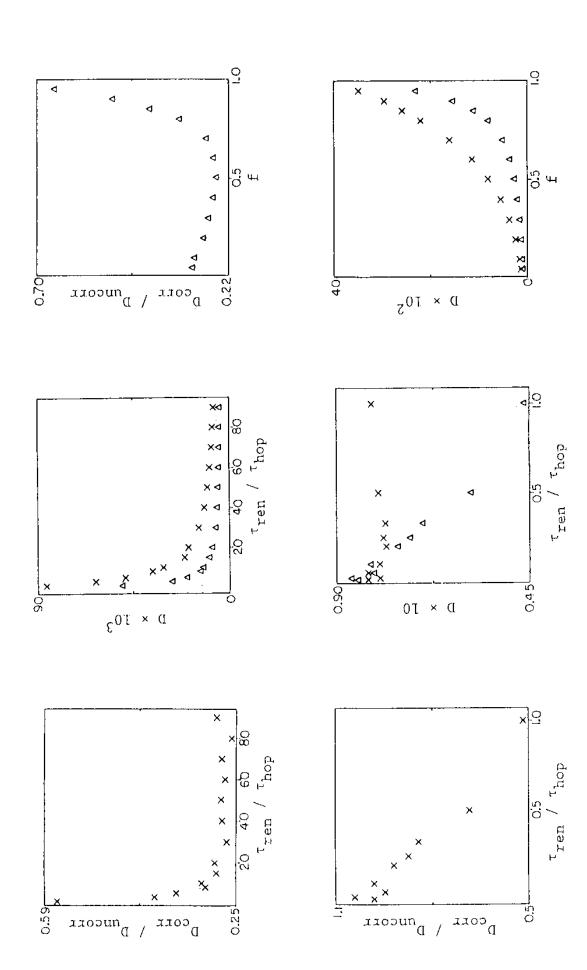
In the dynamic bond percolation models a hopper is propagated from site to site within a lattice. The probability of the hopper moving from one site to another across a bond, wij, is assigned either the value 0 (bond blocked) or the value w. A given fraction, f, of the bonds are unblocked. These unblocked bonds are randomly distributed throughout the lattice. Thus the probability P₁ for the hopping ion to be located on site i is given by

$$\frac{dP_i}{dt} = \sum_{j} \{w_{ji}P_j - w_{ij}P_i\}, \qquad (1)$$

where the sum runs only over nearest neighbors, and

$$w_{ij} = \begin{cases} 0 & \text{probability (1-f)} \\ w & \text{probability f} \end{cases}$$
 (2)

Equation (1) is the kinetic equation, or master equation, for the system. In static bond percolation theory, the bond w_{ij} for any given pair of sites i and j is fixed once and forever to be either available $(w_{11} = w)$ or closed $(w_{ij} = 0)$. In dynamic percolation models, the assignment of wil to either zero or w varies with time. Thus, to add dynamics to the system the assignment of blocked bonds is randomly redistributed, keeping the overall fraction f constant, on a time scale, τ_{ren} . The other relevant parameters of the model are the observation and hopping times, τ_{obs} and τ_{hop} , respectively. Numerical simulations and a closed-form evaluation 10,11 for a onedimensional model show that diffusion occurs in the limit of $\tau_{obs} > \tau_{ren}, \tau_{hop}$. When no fluctuations of the lattice are allowed to occur, $\tau_{ren} > \tau_{obs}$, diffusion is prohibited below the percolation threshold, 10 This second limit, $\tau_{ren} > \tau_{obs}$, corresponds to a polymer electrolyte that is either highly crystalline or is below its glass transition temperature with a resulting low conductivity.



one dimension with correlated and uncor-(b) The individual diffusion increases from 0.05 to 0.95; Δ , D_{corr} ; (a) The ratio of the diffusion coeffi-Dcorr/Duncorr, as a clents obtained from the DBP model in coefficients both increase as f function of the fraction of related renewals, bonds. Thop increases from I to 90 both The diffusion coefficients obtained from (g) (b) As Δ, correlated

function of the relative renewing and

the DBP model in one dimension as a

renewals; X, uncorrelated renewals.

hopping times, tren/thop.

diffusion coefficients decrease.

ren'thop

tren/Thop

decreases from 1 to 0.01 (f = 0.20)

As Tren/

renewal and hopping

function of relative

lated renewals,

(a) Dcorr/Duncorr

decreases to a constant value as the

thop.

times, Tren/

tren/thop increases from 1

The ratio of the diffusion coefficients

dimension with correlated and uncorre-

X, Duncorr (Tren/Thop = 20).

Increases while Duncorr, X,

remains constant (f = 0.20).

To make the model more physically reasonable, correlated instead of $random^{9-11}$ bond renewals are used. This better represents the important segmental fluctuations of the polymer host in polymer electrolytes. To render bond renewal a strongly correlated process a restriction is added such that a bond can renew only by exchanging blocked/unblocked status $(w_{ij} = 0 \text{ or } w_{ij} = w) \text{ with one of its nearest}$ neighbors. This exchange occurs with a characteristic timescale, tren. The results of numerical simulations on a one-dimensional lattice give the mean-squared displacement of the hopper as a linear function of time, the slope being the diffusion coefficient. The effect of the added limitation of correlated renewal is determined by examining the relative diffusion coefficients of the correlated and uncorrelated systems, $D_{\rm C}/D_{\rm u}$, as a function of the parameters, f, Tren, Thop.

Correlations in the renewal process reduce the diffusion coefficient such that $D_c/D_u < 1$, Figures 1-3. At two extreme limits, f + 1.00 τ_{ren} , the value of D_c/D_u approaches unity, Figures 1b and 3a. For the case of $f \rightarrow 1.00$, the renewal process becomes unimportant since there are essentially no blocking bonds. For the case of $\tau_{\mbox{\scriptsize hop}}$ >> $\tau_{\mbox{\scriptsize ren}}$ the lattice fluctuates rapidly relative to the hopper motion. In the time between one hop and the next the lattice passes through such a large number of configurations that all correlations are lost and the correlated renewals become effectively random. The value of the ratio Dc/Du decreases as $\tau_{\text{ren}}/\tau_{\text{hop}}$ increases, illustrated in Figure 1. For $\tau_{ren} > \tau_{hop}$ the hopper samples a cluster of connected sites in the time between consecutive renewals. During this time the hopper encounters blocked bonds which bar its movement. Uncorrelated renewals can open up new areas of available sites while correlated renewals only free one site per renewal, limiting diffusion.

As τ_{ren} becomes large the hopper samples all sites available to it before the next renewal occurs and the effects of correlation reach a limit. This limit is observed as D_c/D_u reaches a constant value, Figure la.

The dependence of $D_{\rm C}/D_{\rm u}$ on the bond fraction f reflects the dynamics of the bond diffusion itself. Like any system of diffusing entities with short range (e.g., excluded volume) interactions this diffusion process is slowest at f=0.5, and roughly proportional to f(1-f). Accordingly the factor $D_{\rm C}/D_{\rm u}$ goes through a similar minimum as shown in Figure 3a. The lack of symmetry of the $D_{\rm C}/D_{\rm u}$ vs. f curve results from the fact that $D_{\rm C}$ is related to the diffusion of the ions which obviously increases strongly with f.

The independent behavior of the diffusion coefficients for the cases of correlated and uncorrelated bond renewal as a function of f and $\tau_{\rm ren}/\tau_{\rm hop}$ are plotted in Figures 2 and 3b. Relating the system to polymer electrolytes, an increase in f and decrease in $\tau_{\rm ren}/\tau_{\rm hop}$ corresponds to an increase in free volume and configurational entropy and a decrease in percent crystallinity of the polymer host, resulting in an improvement of ionic conductivity. These factors are true of both correlated and uncorrelated $^{9-11}$ bond renewals.

An important feature of the DBP model is that diffusion can still occur when the fraction of unblocked bonds is below that which defines the percolation threshold in static percolation models! (f < 1.0 for one dimension). Extending the simulations to two dimensions, it will be of considerable interest to see if there are any sudden changes in the behavior of the diffusion coefficient at the percolation thresholds found in the static models. Our work as well as that of Harrison and Zwanzig! in models with independent (uncorrelated) renewals shows that although

there is substantial change in the quantitative behavior of D(f) as f passes through the threshold value f_c , the abrupt change of the static model is very much smoothed by renewal. We expect this general behavior to occur also for strongly correlated renewal and beyond the effective medium approximation. Experimentally, this means that independent of carrier concentration or of polymer chain length, diffusion of mobile ions will be expected above the glass transition temperature.

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